



SIMULATION AND MODELLING TURBULENT SPRAY DYNAMICS

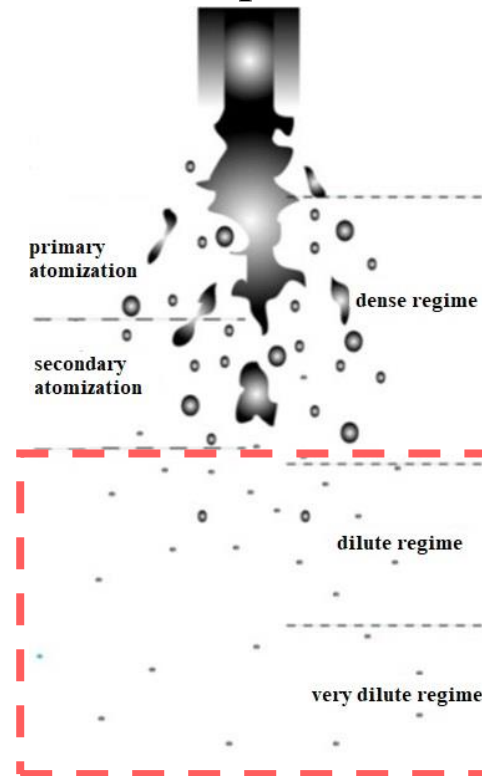
DOCTOR STUDENT: **Jietuo WANG**
SUPERVISOR: **Prof. Francesco Picano**
DATE: 11/09/2020

□ Turbulent spray

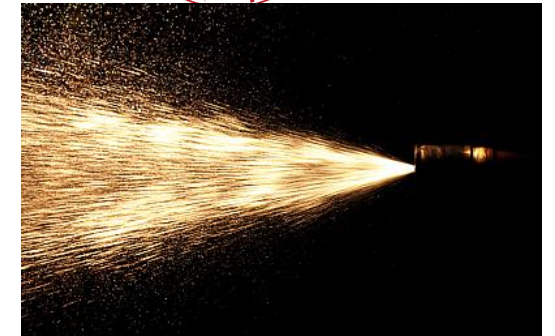
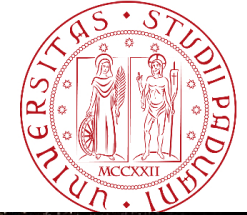
- Complex multiphase flows where two distinguished phases mutually interact exchanging mass, momentum and energy in a turbulent environment.
- Evaporation, preferential concentration, mixing process, combustion, pollutant formation.....
- Unsteady, multiphase, multiscale...
- A satisfied understanding? The capability to modeling?

□ Diluted regime

- Dispersed droplets: **Point-droplet approximation**
- **No break-up** : surface tension \gg aerodynamic forces
- No collision / coalescence : **low volume fraction** ($\Phi < 10^{-3}$)
- 1-way or **2-way** coupling method
- Main region occurring evaporation and combustion



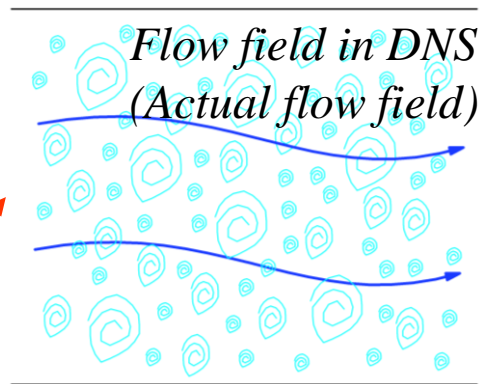
Sketch of various regimes in turbulent sprays [Jenny.2012]



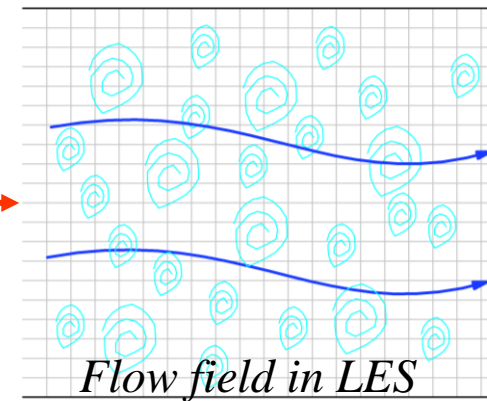
Description of numerical approaches for turbulent flows



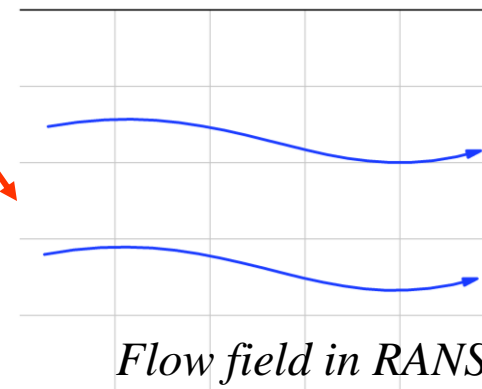
❑ Basic tool for research



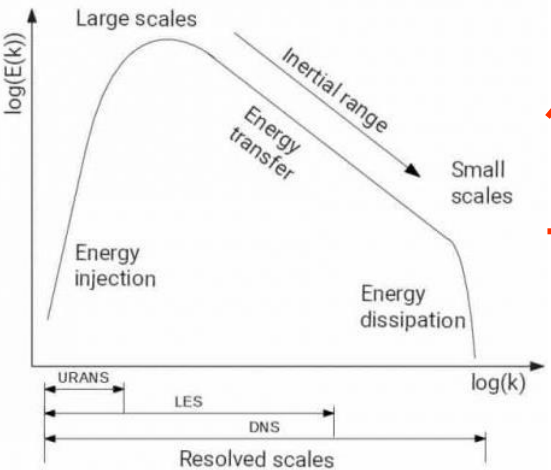
- More flow details available



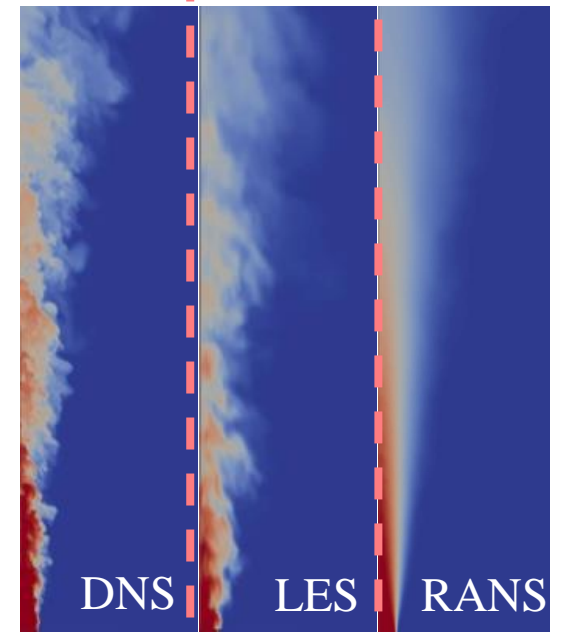
- Decent computational cost



❑ Popular tools for industrial application



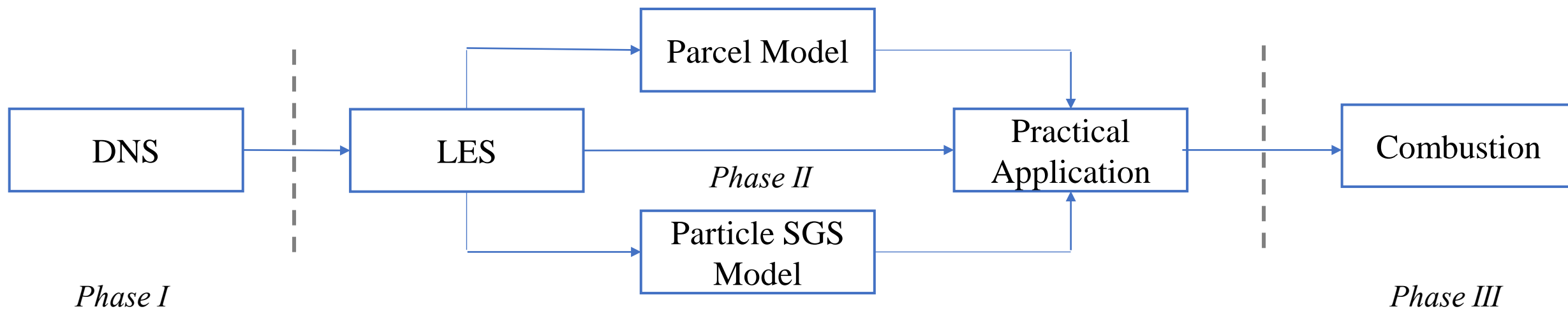
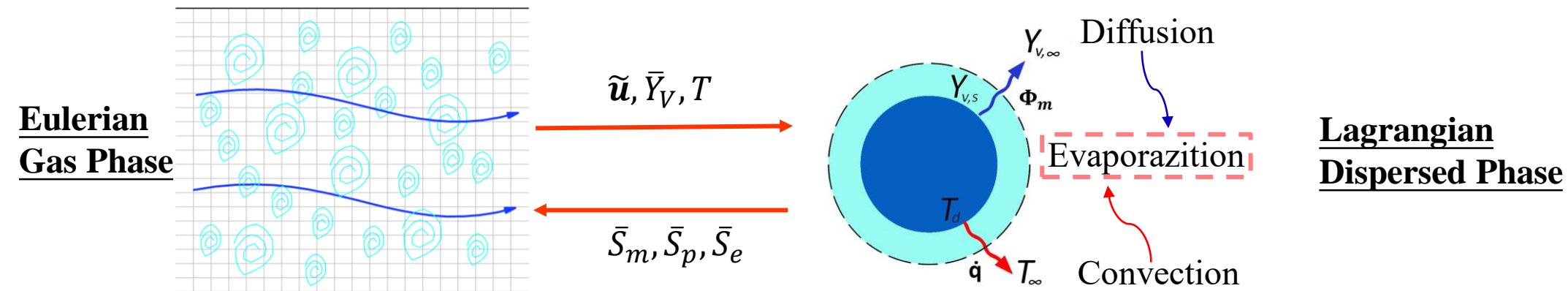
Increase in resolution & computational cost



Comparison between numerical methods for flow

The question & our answer

- ❑ Is it possible to further enhance the computational efficiency in LES approach?
- ❑ Is it possible to further improve the prediction accuracy in LES approach?



Numerical tool and Formulations for DNS

□ Numerical Tool: CYCLONE

- Fully turbulent flow
- Staggered mesh
- Low Mach number NS & Point-droplet equations
- Cylindrical coordinate
- MPI parallelization



□ Formulations for DNS

Eulerian Gas Phase
(Low Mach Navier-Stokes)

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = S_m$$

$$\frac{\partial}{\partial t} (\rho Y_V) + \nabla \cdot (\rho Y_V \mathbf{u}) = \nabla \cdot (\rho \mathcal{D} \nabla Y_V) + S_m$$

$$\frac{\partial}{\partial t} (\rho \mathbf{u}) + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u}) = \nabla \cdot \boldsymbol{\sigma} - \nabla P + S_p$$

$$\frac{\partial}{\partial t} (\rho E) + \nabla \cdot (\rho E \mathbf{u}) = -\nabla P \mathbf{u} + \nabla \cdot (\boldsymbol{\sigma} \otimes \mathbf{u}) - \nabla q + S_e$$

Lagrangian dispersed phase
(Point-droplet equations)

$$\frac{d\mathbf{x}_d}{dt} = \mathbf{u}_d$$

$$\frac{d\mathbf{u}_d}{dt} = \frac{(\mathbf{u} - \mathbf{u}_d)}{\tau_d} (1 + 0.15 \text{Re}_d^{0.687})$$

$$\frac{dr_d^2}{dt} = -\frac{\mu_g}{\rho_l} \frac{\text{Sh}}{\text{Sc}} \ln(1 + B_m)$$

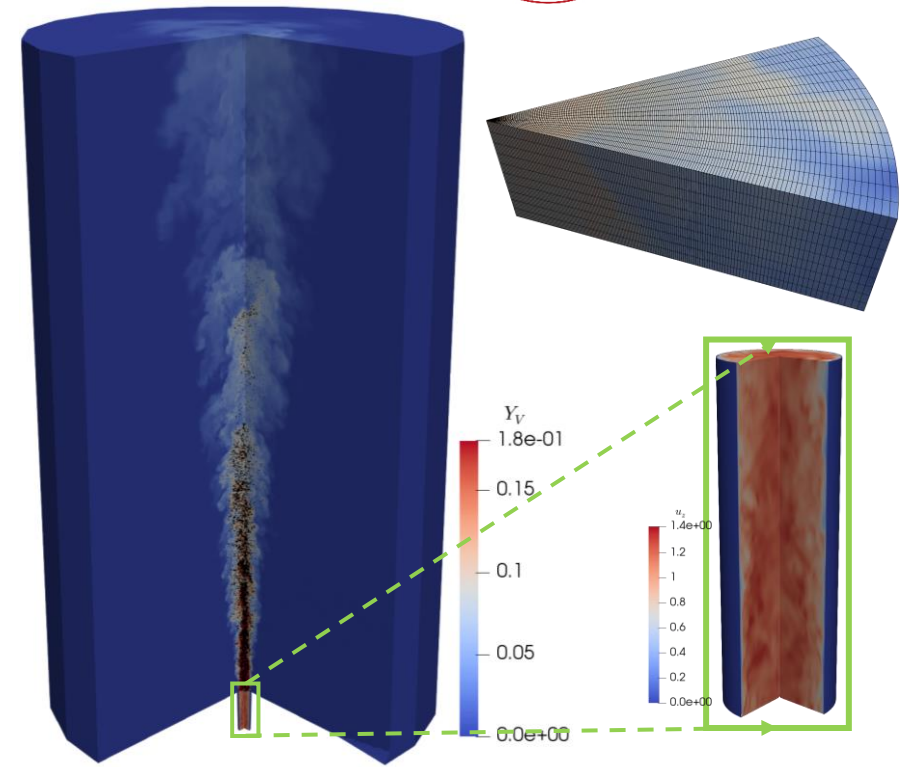
$$\frac{dT_d}{dt} = \frac{\text{Nu}}{3\text{Pr}} \frac{c_{p,g}}{c_{p,l}} \frac{T - T_d}{\tau_d} + \frac{L_v}{c_{p,l}} \frac{\dot{m}_d}{m_d}$$

2-way coupling terms

$$S_m = -\sum_{i=1} \frac{dm_{d,i}}{dt} \delta(\mathbf{x} - \mathbf{x}_{d,i})$$

$$S_e = -\sum_{i=1} \frac{d}{dt} (m_{d,i} c_{p,l} T_{d,i}) \delta(\mathbf{x} - \mathbf{x}_{d,i})$$

$$S_p = -\sum_{i=1} \frac{d}{dt} (m_{d,i} \mathbf{u}_{d,i}) \delta(\mathbf{x} - \mathbf{x}_{d,i})$$

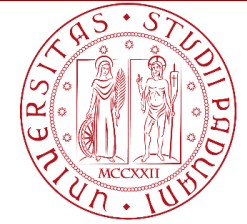


A sketch of the 3D cylindrical domain

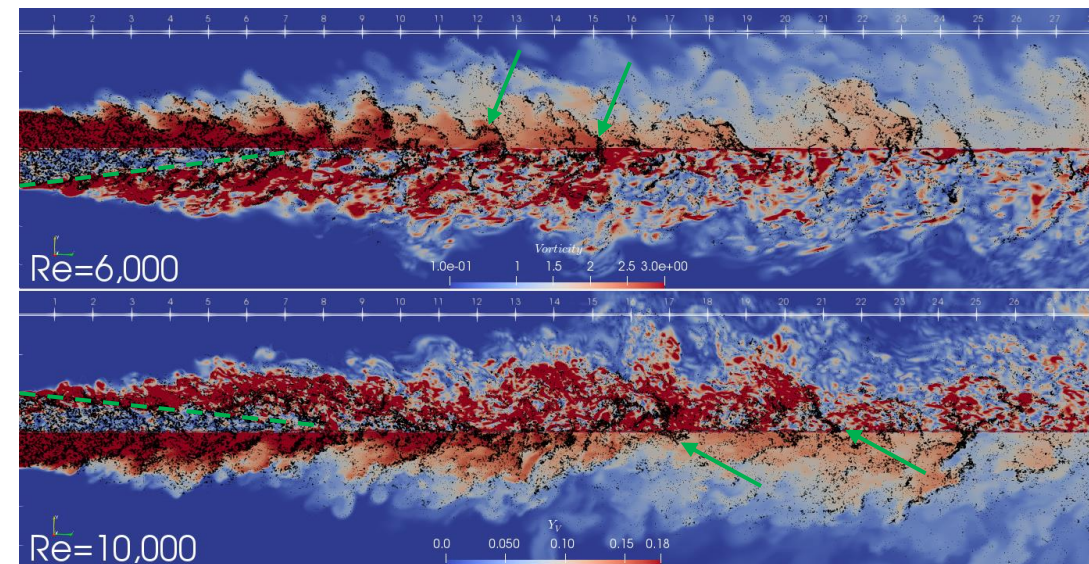
More details: Battista et al. PoF 2011 and Dalla Barba & Picano PRF 2018

Benchmark simulation parameters

- ❑ Benchmark simulation parameters (*DNS*)
 - Monodisperse acetone droplets at the inflow ($r_{d,0}=6\mu\text{m}$)
 - Turbulent inflow with saturated gas ($S=0.99$, $T=275\text{ K}$)
 - Reynolds number : $\text{Re}=2U_0R/\nu=6,000$ & $10,000$ ($U_0=8.1\text{ m/s}$ & $U_0=13.9\text{ m/s}$)
 - Quiescent environment of dry air
 - Non-uniform mesh 46 M points
 - ~3M evaporating droplets with mass fraction $\Phi\approx 0.05$

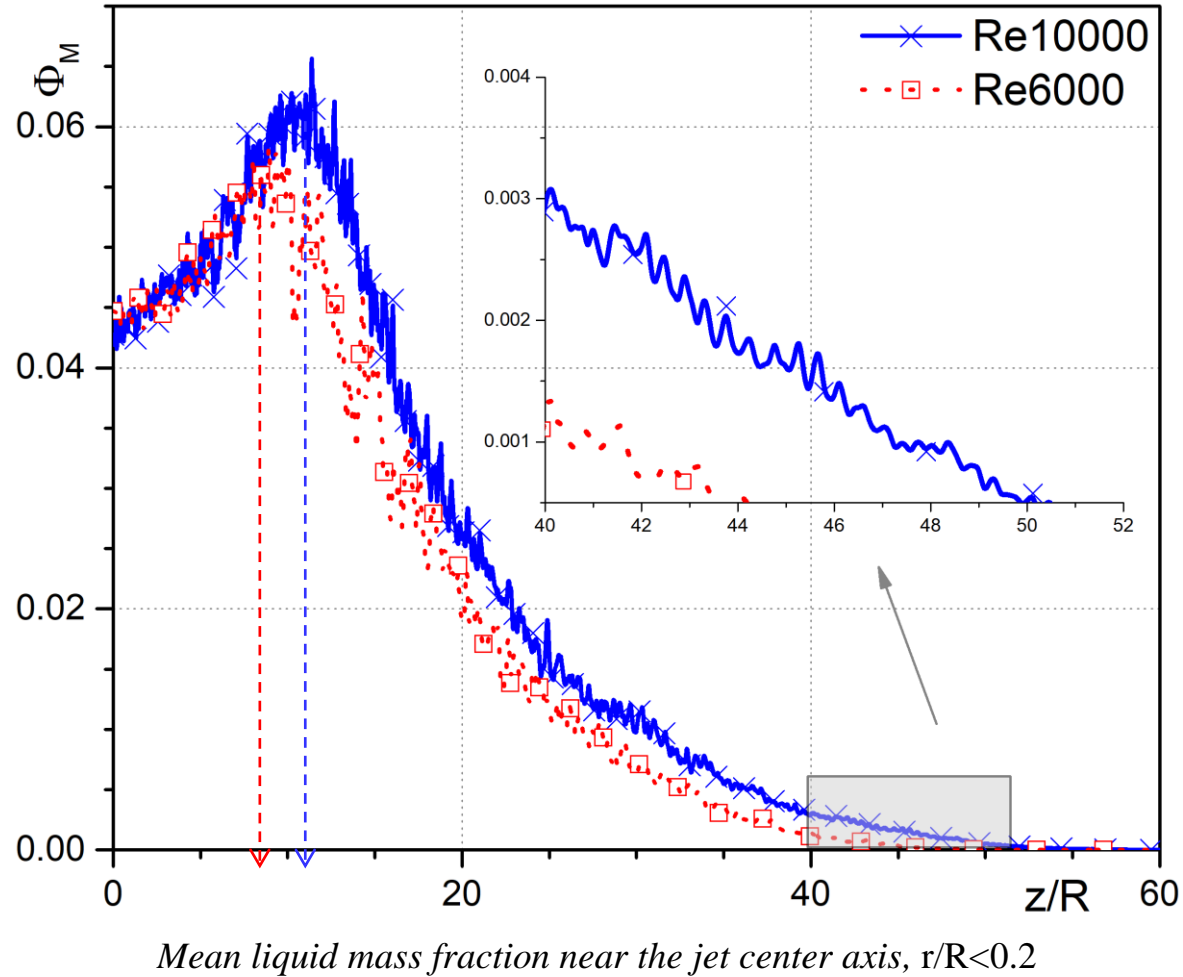
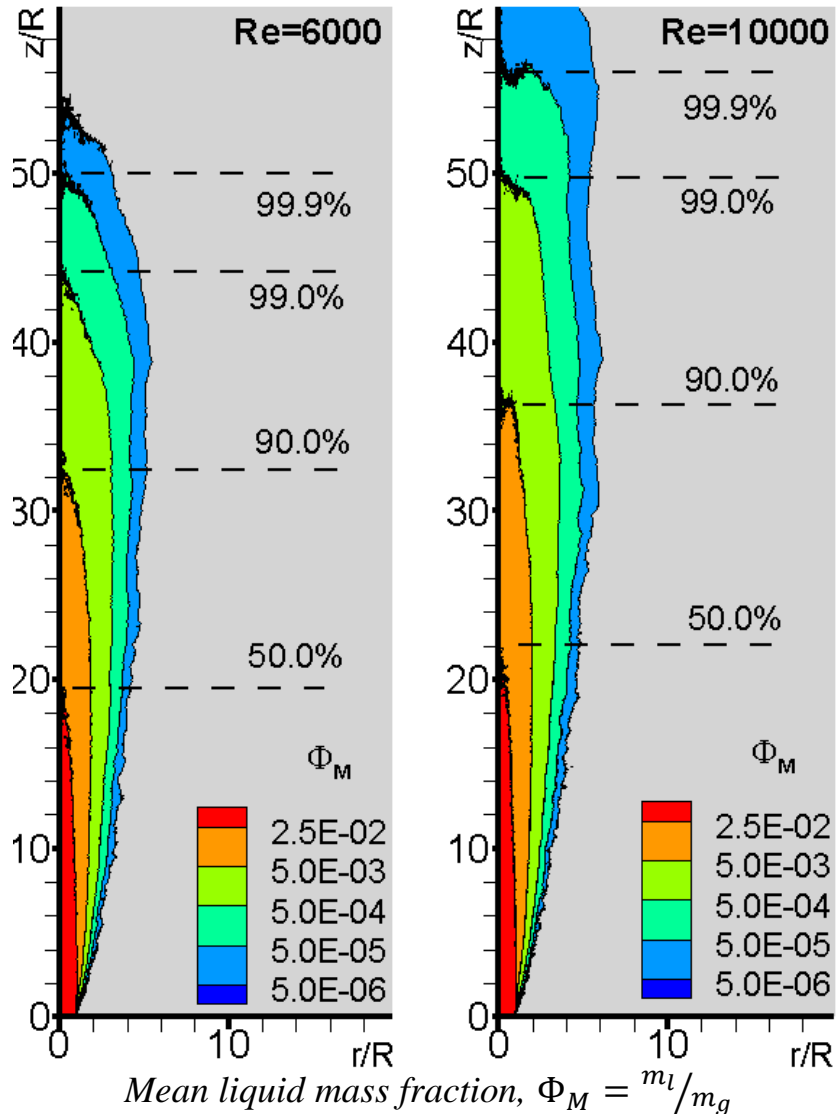


Radial slices of instantaneous distribution for vapor mass fraction, $Y_v = \rho_v/\rho_g$, and vorticity.



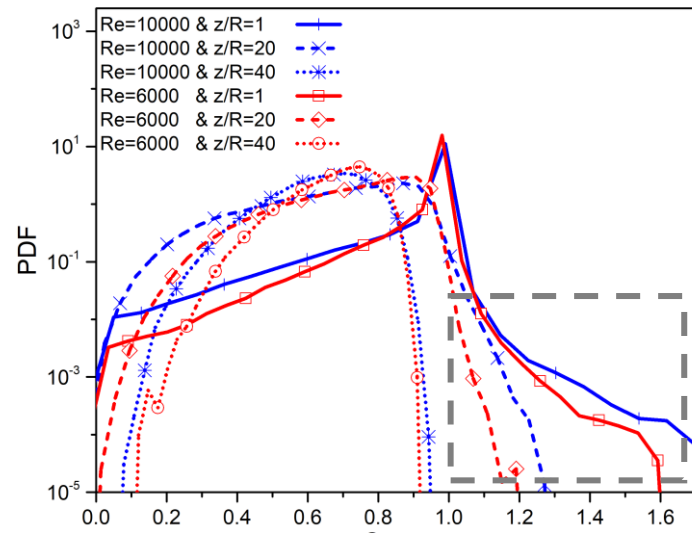
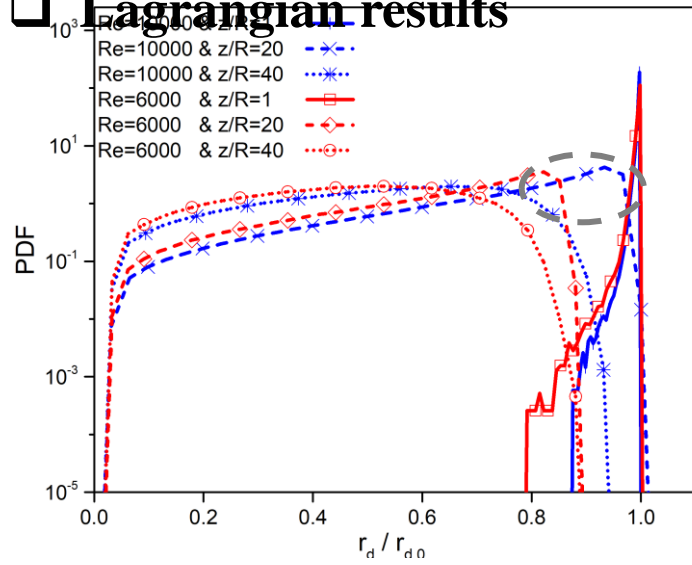
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- ❑ A qualitative view
 - The preferential concentration is evident for both cases.
 - A longer self-potential core is observable in $\text{Re}=10,000$ case.

□ Mean field results



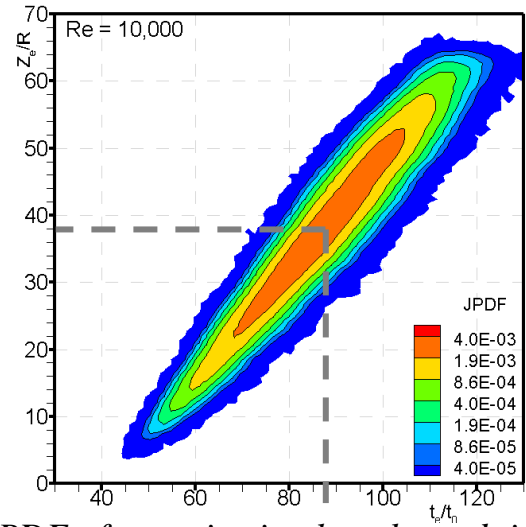
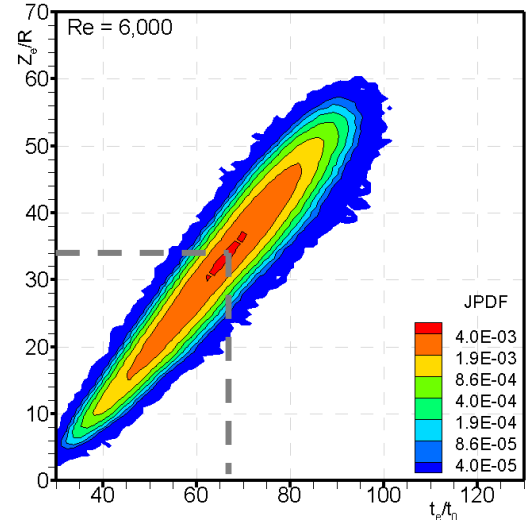
- A longer vaporization length (99%) is shown in Re = 10,000 case.
- The Φ_M exhibits a near-field hump shape along the center axis.

□ Lagrangian results



PDF of non-dimensional droplet radius and the saturation field at droplets surface

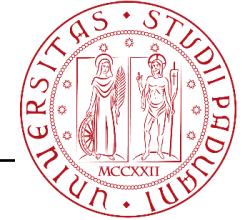
- Towards the intermediate-fields and far-fields region, **more robust spread of the droplet size spectrum is shown in $Re = 10,000$ case.**
- Certain amount of droplets stay in range $S_d > 1$, i.e. supersaturation, which is more intensive in higher Re sprays.



JPDF of vaporization lengths and times of the droplets computed over the whole droplet population



- A strong linear correlation between the droplet vaporization length and time while different slope rates are observable.
- In $Re = 10,000$ case, the median vaporization times is $t_e/t_0 = 90$, which is larger than its correspondent value, $t_e/t_0 = 70$ in $Re = 6,000$ case.



Eulerian Gas Phase (Low Mach Navier- Stokes)

$$\frac{\partial \bar{\rho}}{\partial t} + \nabla \cdot (\bar{\rho} \tilde{\mathbf{u}}) = \bar{S}_m$$

$$\frac{\partial}{\partial t} (\bar{\rho} \tilde{Y}_V) + \nabla \cdot (\bar{\rho} \tilde{Y}_V \tilde{\mathbf{u}}) = \nabla \cdot (\bar{\rho} \tilde{D} \nabla \tilde{Y}_V) - \nabla q_{Y_V} + \bar{S}_m$$

$$\frac{\partial}{\partial t} (\bar{\rho} \tilde{\mathbf{u}}) + \nabla \cdot (\bar{\rho} \tilde{\mathbf{u}} \otimes \tilde{\mathbf{u}}) = \nabla \cdot \bar{\boldsymbol{\sigma}} - \nabla \cdot \boldsymbol{\tau}^R - \nabla \bar{P} + \bar{S}_p$$

$$\frac{\partial}{\partial t} (\bar{\rho} \tilde{e}) + \nabla \cdot (\bar{\rho} \tilde{e} \tilde{\mathbf{u}}) = -\nabla \bar{P} \tilde{\mathbf{u}} + \nabla \cdot (\bar{\boldsymbol{\sigma}} \otimes \tilde{\mathbf{u}}) - \nabla \bar{q} - \nabla q_e + \bar{S}_e$$

Lagrangian dispersed phase (Point-droplet equations)

$$\frac{d\mathbf{x}_d}{dt} = \mathbf{u}_d$$

$$\frac{d\mathbf{u}_d}{dt} = \frac{(\tilde{\mathbf{u}} - \mathbf{u}_d)}{\tau_d} (1 + 0.15 \text{Re}_d^{0.687})$$

$$\frac{dr_d^2}{dt} = -\frac{\bar{\mu}_g}{\rho_l} \frac{\text{Sh}}{\text{Sc}} \ln(1 + B_m)$$

$$\frac{dT_d}{dt} = \frac{\text{Nu}}{3\text{Pr}} \frac{\bar{c}_{p,g}}{c_{p,l}} \frac{T - T_d}{\tau_d} + \frac{L_v}{c_{p,l}} \frac{\dot{m}_d}{m_d}$$

$$\bar{S}_m = \int S_m G_\Delta(\mathbf{x}, \mathbf{r}) d\mathbf{r}$$

$$\bar{S}_e = \int S_e G_\Delta(\mathbf{x}, \mathbf{r}) d\mathbf{r}$$

$$\bar{S}_p = \int S_p G_\Delta(\mathbf{x}, \mathbf{r}) d\mathbf{r}$$

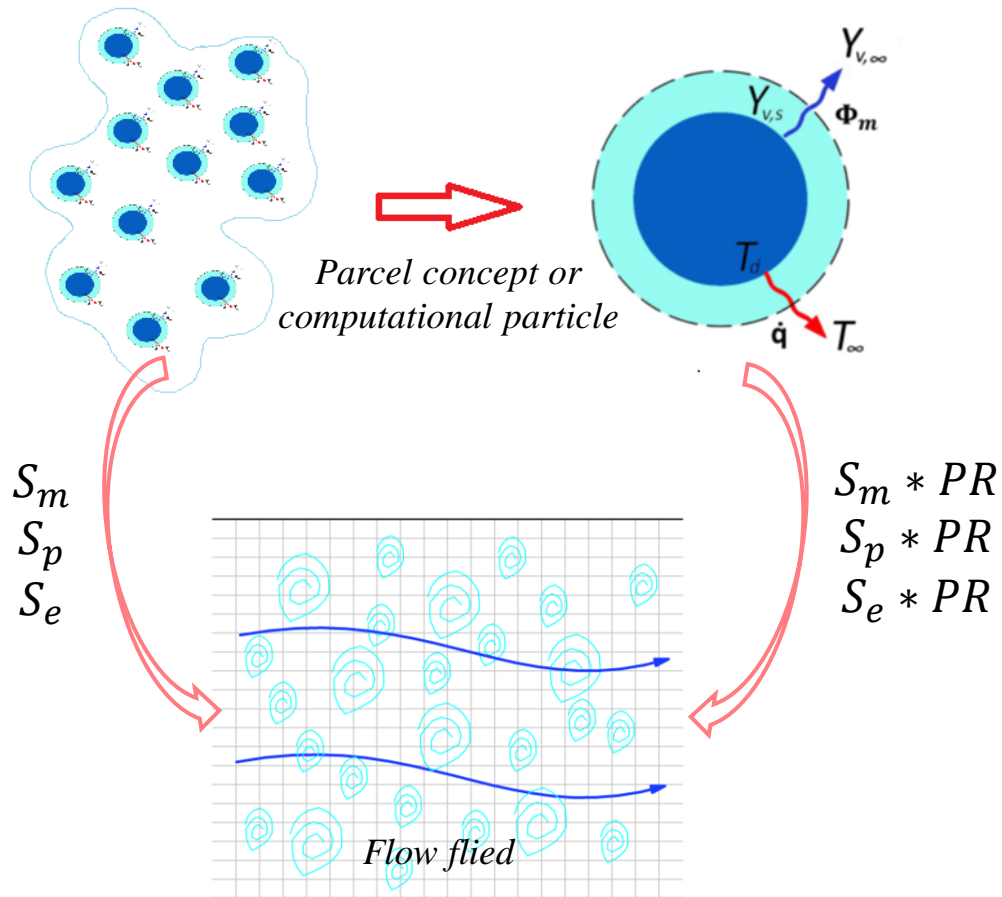
Boussinesq Hypothesis: $\boldsymbol{\tau}^R = -2\mu_{SGS}(\tilde{S} - 1/3\tilde{S}\mathbf{I}) \quad \tilde{S} = 0.5(\nabla \tilde{\mathbf{u}} + \nabla \tilde{\mathbf{u}}^T)$

Smagorinsky Model: $\mu_{SGS} = \bar{\rho}(C_S \Delta)^2 |\tilde{S}| \quad |\tilde{S}| = (2\tilde{S}\tilde{S})^{1/2} \quad \Delta = (r\Delta\theta\Delta r\Delta z)^{1/3}$

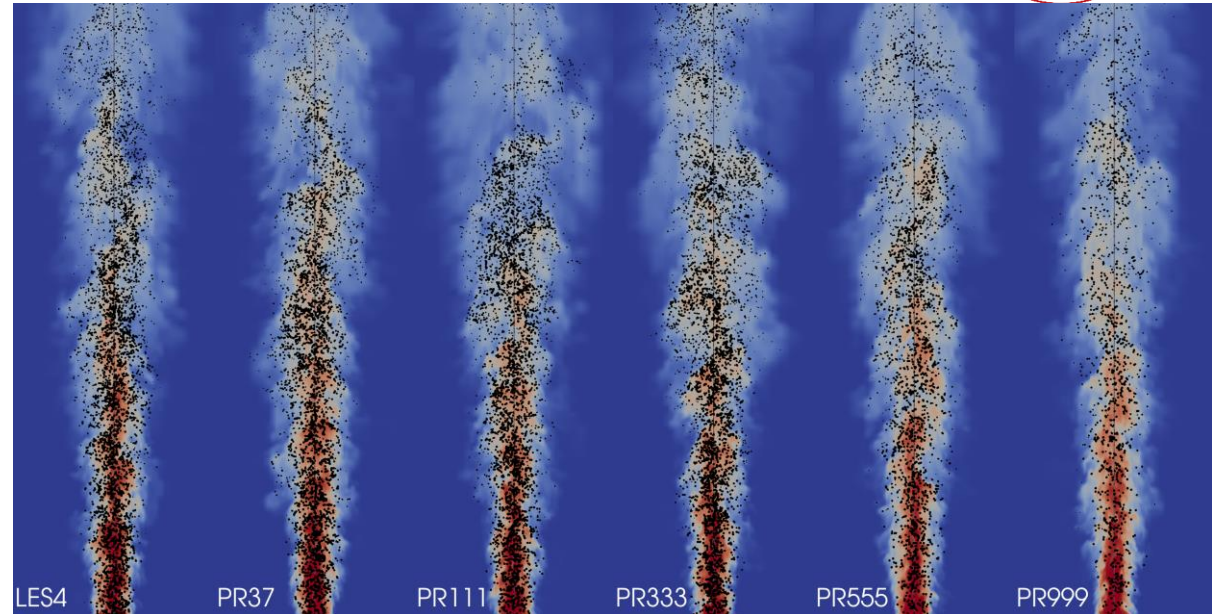
- The mesh for LES is coarsen from DNS mesh by increasing mesh size 4 times in each direction. The total number is ~ 0.7 M points.

❑ The classical Smagorinsky SGS model is employed in the LES approach.

❑ All LES simulations are based on formulations listed above.



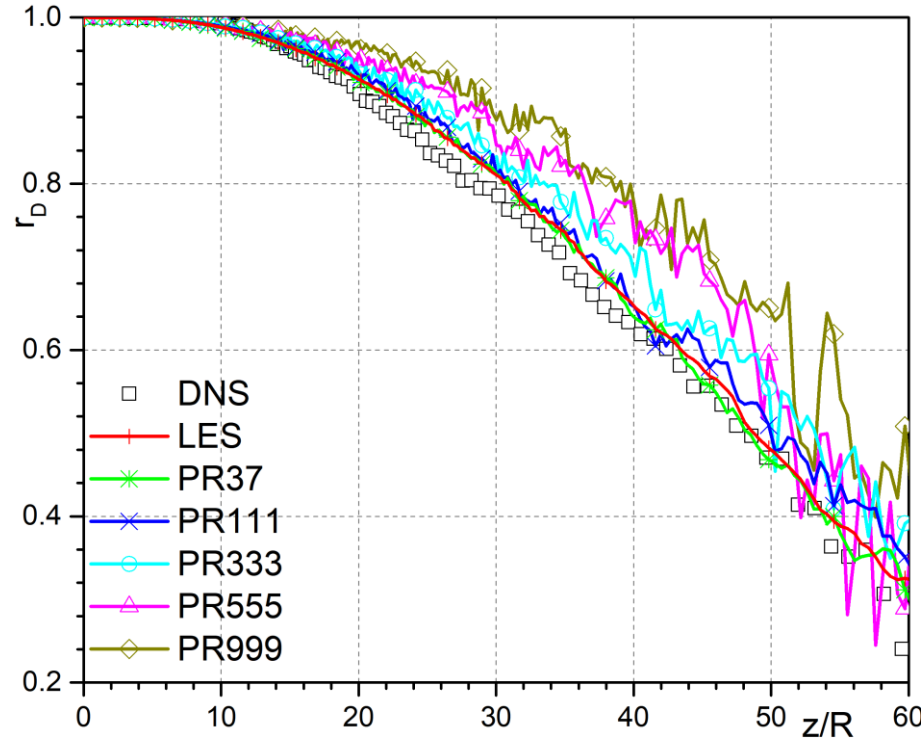
Results



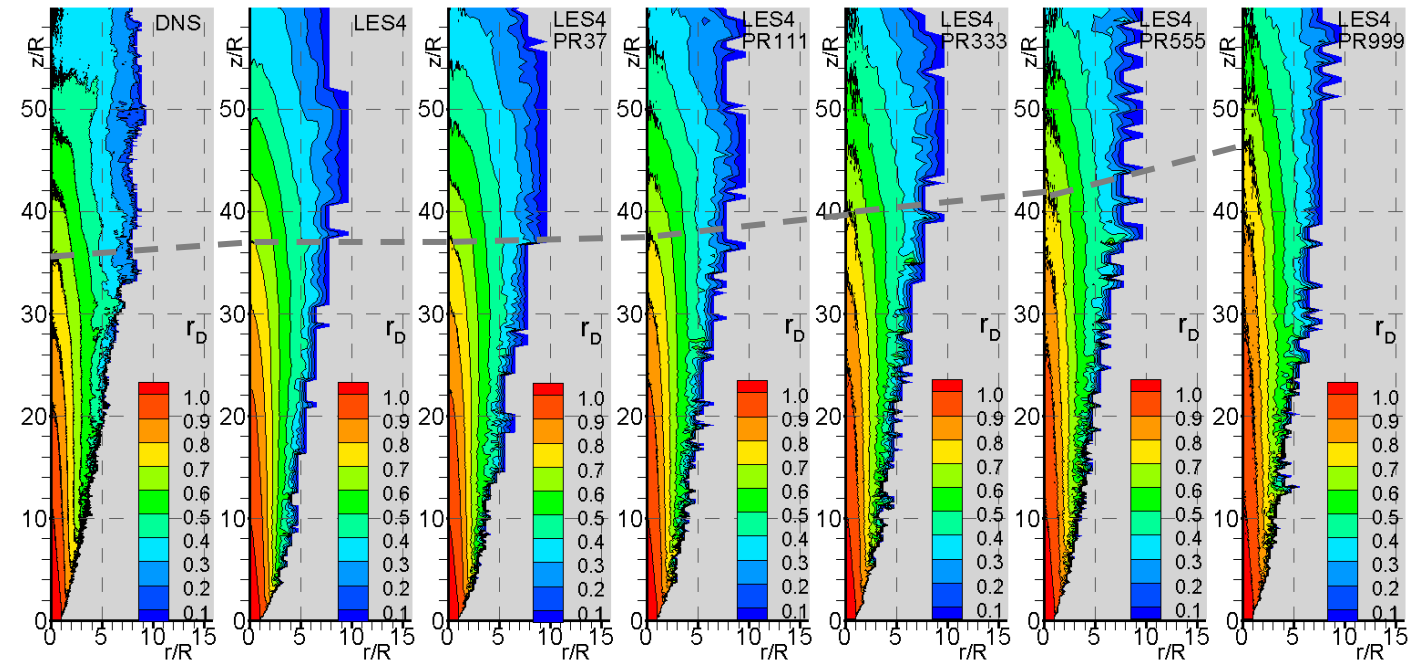
Radial slices of instantaneous distribution for vapor mass fraction, Y_v , with different Parcel Ratios (PR)

- A cluster of droplets with similar properties are represented by one “parcel”, e.g. temperature, radius and velocity.
- **Parcel model is added in our code package successfully.**
- From the qualitative point-of-view, no remarkable difference is observable among LES simulations with different PRs.

Results



Mean droplet radius near the jet center axis, $r/R < 0.2$



Mean droplet radius rescaled by the droplet initial radius, r_d , with different Parcel Ratios (PR)

- Droplet vaporization: The lifetime of droplets could be prolonged if large PR values are used in modeling spray dynamics.
- **PR = 111** is recommended considering the computational efficiency and prediction accuracy.

Model description

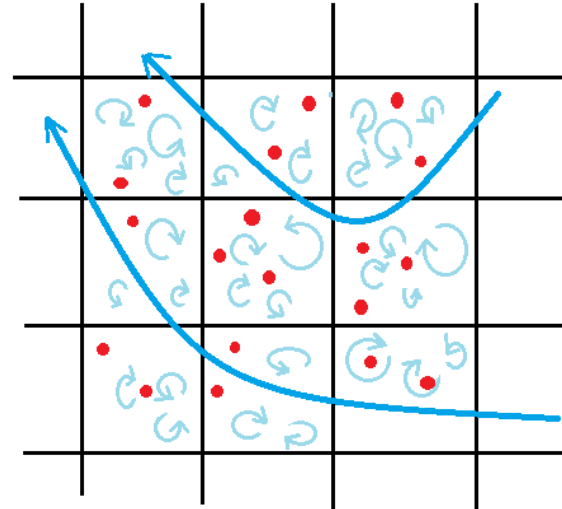
$$\frac{d\mathbf{u}_d}{dt} = \frac{(\tilde{\mathbf{u}} + \mathbf{u}_{sgs} - \mathbf{u}_d)}{\tau_d} (1 + 0.15\text{Re}_d^{0.687})$$

$$\mathbf{u}_{sgs} = \chi \sqrt{\frac{2}{3} \frac{\mu_t}{\rho C_k \Delta \text{Re}}}, \text{ (SGS velocity)}$$

$$\chi \sim \mathcal{N}(0,1), \text{ (Gaussian distribution)}$$

$$\tau_{sgs} = \frac{\bar{\Delta} \text{Re}}{\mu_t} \ \& \ \tau_{sgs} = \tau_d^{1.6} / \left(\frac{\bar{\Delta} \text{Re}}{\mu_t} \right)^{0.6}$$

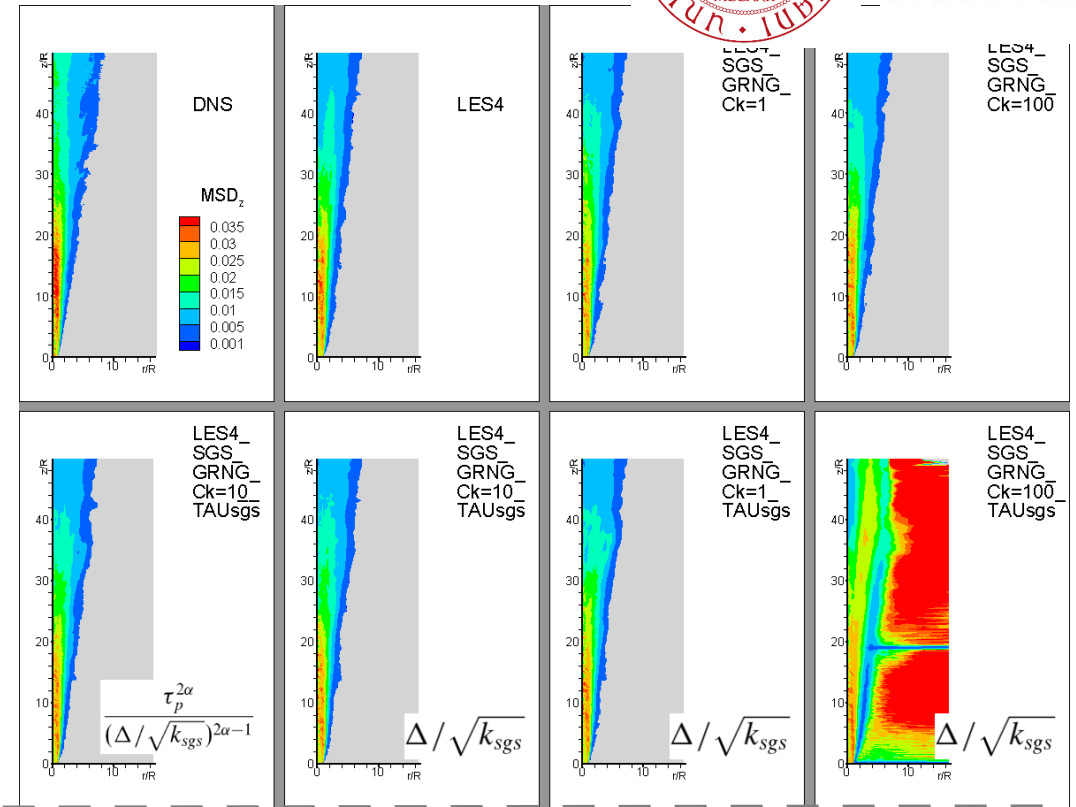
, (SGS time scale)



Top: Schematics of particles in LES framework

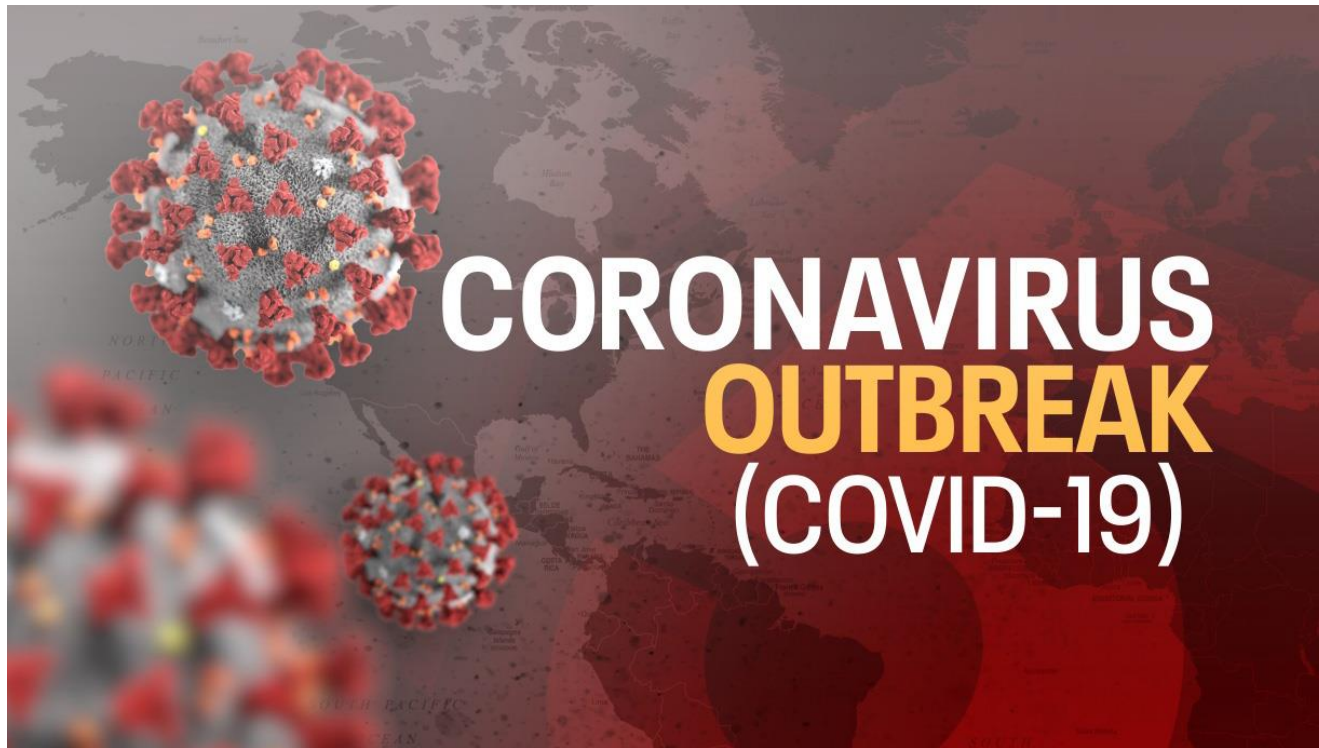
Right: Mean square displacement (MSD) in z direction for DNS, original LES and LES cases with different particle SGS models or model coefficients

Results



- A stochastic model to account for the SGS part in particle momentum equation is proposed and implemented in our code.
- The Ziggurat method [Marsaglia.2000] is adopted to generate a Gaussian random number for the SGS model.
- **No significant improvement compared to the original LES has been observed.** Further investigations is undergoing.

□ COVID-19



- To fill scientific gaps in our understanding of critical issues.
- To better prepare ourselves to tackle the next break of COVID-19 or a similar disease.
- **The behaviors of droplets laden with virus under different ambient conditions.**

Small droplet aerosols in poorly ventilated spaces and SARS-CoV-2 transmission

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Focus on Fluids
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The flow physics of COVID-19

Published Online
May 27, 2020
[https://doi.org/10.1016/S2213-2600\(20\)30245-9](https://doi.org/10.1016/S2213-2600(20)30245-9)

Rajat Mittal^{1,2,†}, Rui Ni^{1,†} and

THE NEW ENGLAND JOURNAL OF MEDICINE

CORRESPONDENCE

Aerosol and Surface Stability of SARS-CoV-2 as Compared with SARS-CoV-1

JAMA Insights

**Turbulent Gas Clouds and Respiratory Pathogen Emissions
Potential Implications for Reducing Transmission of COVID-19**

Lydia Bourouiba, PhD

Journal of Aerosol Science

journal homepage: <http://www.elsevier.com/locate/jaerosci>

ORIGINAL ARTICLE

Influence of wind and relative humidity on the social effectiveness to prevent COVID-19 airborne transmission from healthy and influenza-infected subjects: a numerical study

Yu Feng^{a,†}, Thierry Marchal^b, Ted Sperry^a, Hang Yi^a

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Experimental investigation of far-field human cough airflows in obstructing respiratory jets

Nicholas Dudalski¹ | Ahmed Mohamed¹ | Samira Mubareka^{2,3} | Ran Bi¹ | Chao Zhang¹ | Eric Savory¹

The graphical abstract displays the growth of COVID-19 infections worldwide as of April 20, 2020, superimposed on the result from a direct-numerical simulation by Jung-Hee Seo (Johns Hopkins University) and Kourosh Shoole (Florida State University), showing the vortices generated by a cough through a face mask.

Visualizing the effectiveness of face masks in obstructing respiratory jets

Cite as: Phys. Fluids 32, 061708 (2021). <https://doi.org/10.1063/5.0016018> | Submitted: 31 May 2020, Accepted: 06 June 2020, Published Online: 30 June 2020

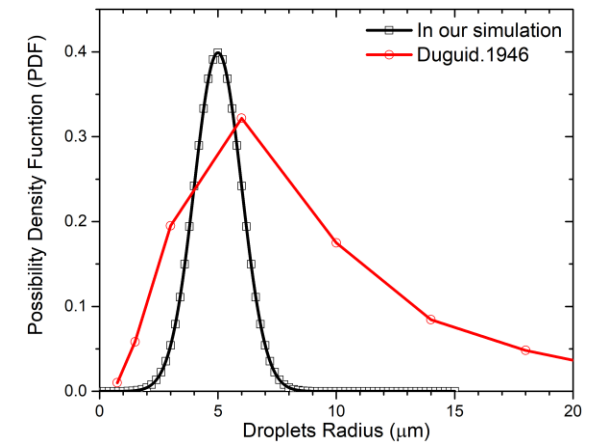
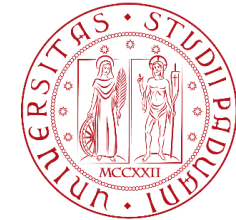
❑ Coughing jets: Initial Conditions, Boundary Conditions and Parameters Setup

Parameters	Value
Inlet radius[m]	1×10^{-2}
Inlet temp.[K]	310.15
Inlet RH[%]	100
Max. inlet vel.[m/s]	10.93
Initial pressure[Pa]	101300
Ambient Temp.[K]	296.15
Ambient RH[%]	19 & 50 & 90
Droplets radius[m]	$N(5 \times 10^{-6}, 1)$
Liquid volume fraction[%]	1.9×10^{-7}
Bulk Re	12000
Computational domain ($\theta \times r \times z$)[m]	$2\pi \times 1 \times 2$
Mesh	$64 \times 280 \times 600$
Time step [s]	1×10^{-5}

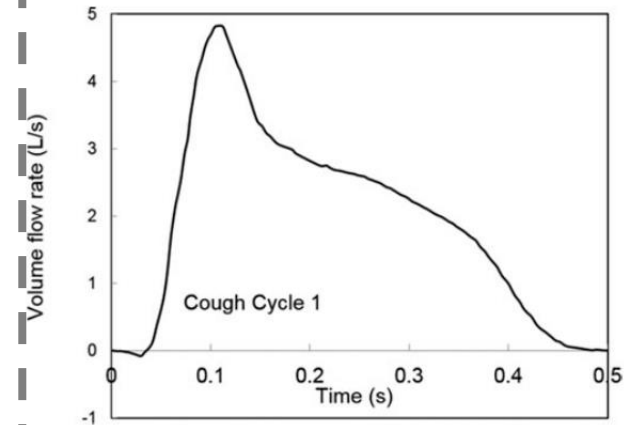


An enhanced image of a sneeze process

- Other thermodynamic parameters for humid air is calculated based on [Picard.2008 & Tsilingiris.2008].
- Buoyancy force is added in both gas phase and particle phase.
- The quantity of injected droplets per time steps is proportional to the prescribed velocity condition.



Droplets radius distribution

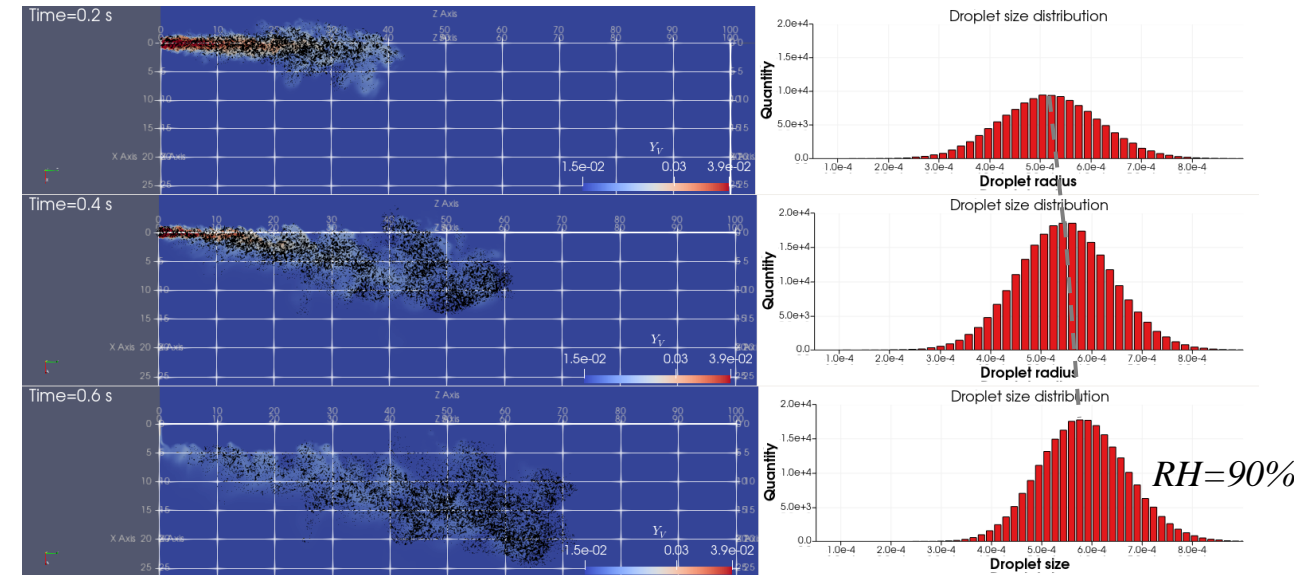
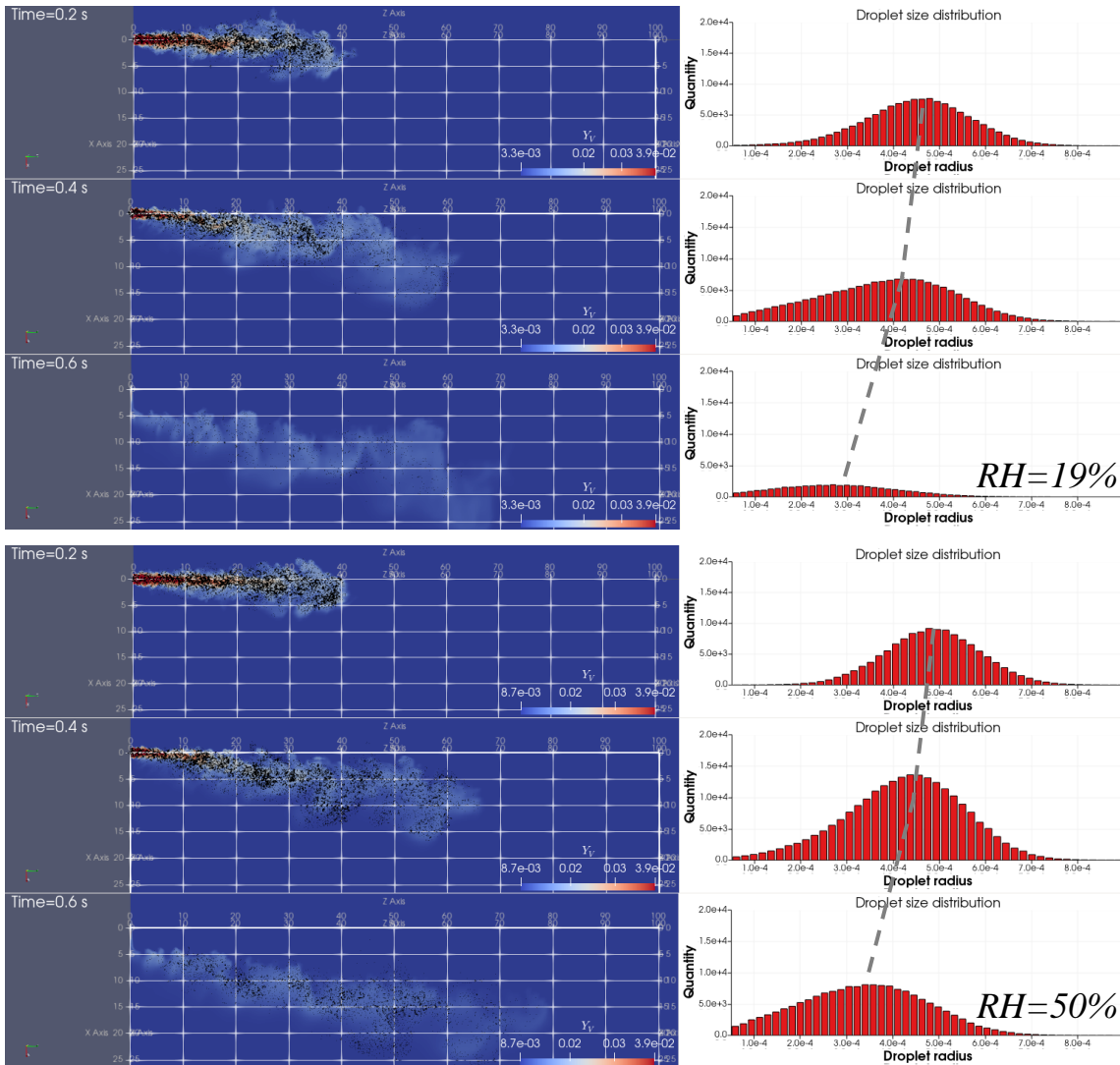


Inlet flow rate [Thatiparti.2016]

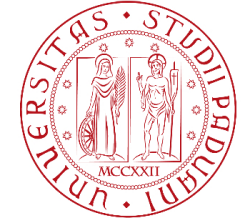
☐ Coughing jets: Preliminary Results



The instantaneous vapor mass fraction distribution sampling with droplets under different Relative Humidity (RH) conditions: 19%, 50%, 90%.



- Slower vaporization is obvious as RH is increased from 19% to 90%.
- **Droplets can “grow up” in high RH conditions.**



- ✓ The spray dynamics of an evaporating turbulent acetone jet in $Re = 10,000$ has been investigated with DNS framework.
 - A longer vaporization length compared to our previous work ($Re=6,000$) was found together with delayed preferential concentration phenomena*.

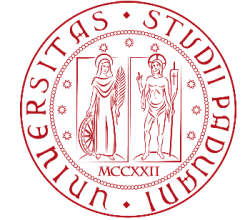
- ✓ 3D LES cases has been performed to obtain the influence of parcel model on Lagrangian phase behaviors.
 - An optimal PR was found to maintain the prediction accuracy*.

- ✓ The impact other submodels on particle dynamics has been checked.

Work in progress

- Particle SGS model
- Coughing jet simulation with different ambient conditions
- A manuscript is preparing for publication*; another one is under writing*.

THE END



Thank you for your attention!